

# Lecture 1 – Some Time-Frequency Transformations

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- What is a Time-Frequency distribution/representation?
- The Wigner Distribution
- The Short-Time Fourier Transform
- The Ambiguity Function
- Uncertainty Principles

# Frequency representation

- Suppose we are given some function  $f \in L^2(\mathbb{R}^d)$ .
- The Fourier transform allows us to realize  $f$  as a superposition of pure frequencies, viz.

$$\widehat{f}(\gamma) = \int_{\mathbb{R}^d} f(x) e^{-2\pi i \langle x, \gamma \rangle} dx, \quad f(x) = \int_{\mathbb{R}^d} \widehat{f}(\gamma) e^{2\pi i \langle x, \gamma \rangle} d\gamma.$$

- Parseval's formula says this realization preserves energy.

$$\int_{\mathbb{R}^d} |f(x)|^2 dx = \int_{\mathbb{R}^d} |\widehat{f}(\gamma)|^2 d\gamma.$$

- The Fourier transform of  $f$  gives a frequency *distribution* of  $f$  as well as a *frequency representation* of  $f$ .

- It is a distribution in the sense that on any set  $S \subseteq \mathbb{R}^d$ , the quantity

$$\int_S |\hat{f}(\gamma)|^2 d\gamma$$

represents the proportion of the function's energy in that frequency band.

- Also the Fourier Inversion gives a representation of  $f$  in terms of its frequency content.
- A mathematically analogous interpretation takes  $x$  as the *position variable* and  $\gamma$  as the *momentum variable*.

# What is a Time-Frequency distribution/representation?

- In the context of acoustic signals  $f(t)$ , it has long been recognized that we do not experience the frequency content of a signal through the usual Fourier paradigm.
- Rather than experiencing a signal as either its *time representation*,  $f(t)$ , or its *frequency representation*,  $\hat{f}(\gamma)$ , a signal consists of different frequencies at different times.
- A simple metaphor for the representation of an acoustic signal jointly in terms of time and frequency is in a musical score (D. Gabor, 1949; de Bruijn, 1965).
- The notion of time-frequency distributions have arisen in the signal processing and quantum physics communities.

Fur Elise  
Clavierstück in A Minor - WoO 59

Ludwig Van Beethoven

Poco meno  
ppp

fl. 1. x

fl.

fl.

fl.

fl.

Public Domain

Fur Elise  
Clavierstück in A Minor - WoO 59

Ludwig Van Beethoven

Poco meno  
ppp

ff

ff

ff

ff

ff

Public Domain



Fur Elise  
Clavierstück in A Minor - WoO 59

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- When a particular frequency (note) is to be played at a particular time, the composer puts a dot at that place on the time (horizontal) and frequency (vertical) axis.



- The musical score can be thought of as
  - the *energy distribution* of the signal  $f$  in time and frequency, and
  - as a *representation* of  $f$  in time and frequency variables.
- Adding the square of the intensities of each note in a certain time and frequency range gives the energy in the signal in that region of time-frequency space.
- At the same time, the signal can be reproduced from the musical score.

- We seek to generalize the idea of the musical score (L. Cohen, 1966, 1989).
- For a given  $f(x)$ , we seek  $P(x, \gamma)$ , a joint distribution giving the intensity of  $f$  at time  $x$  and frequency  $\gamma$ .
- What properties do we require that  $P$  satisfy?

- We require that  $P$  satisfy the *marginals*:

$$\int_{-\infty}^{\infty} P(x, \gamma) d\gamma = |f(x)|^2 \text{ and } \int_{-\infty}^{\infty} P(x, \gamma) dx = |\hat{f}(\gamma)|^2.$$

- The energy of  $f$  in a given time range  $[t_0 - \epsilon, t_0 + \epsilon]$  is

$$\int_{t_0 - \epsilon}^{t_0 + \epsilon} \int_{-\infty}^{\infty} P(x, \gamma) d\gamma dx = \int_{t_0 - \epsilon}^{t_0 + \epsilon} |f(x)|^2 dx.$$

- The energy of  $f$  in a given frequency range  $[\gamma_0 - \delta, \gamma_0 + \delta]$  is

$$\int_{\gamma_0 - \delta}^{\gamma_0 + \delta} \int_{-\infty}^{\infty} P(x, \gamma) dx d\gamma = \int_{\gamma_0 - \delta}^{\gamma_0 + \delta} |\hat{f}(\gamma)|^2 d\gamma.$$

- And obviously,

$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} P(x, \gamma) d\gamma dx = \|f\|_2^2.$$

- We would further like to say that the energy of  $f$  in the time-frequency region

$$[t_0 - \epsilon, t_0 + \epsilon] \times [\gamma_0 - \delta, \gamma_0 + \delta]$$

is

$$\int_{t_0 - \epsilon}^{t_0 + \epsilon} \int_{\gamma_0 - \delta}^{\gamma_0 + \delta} P(x, \gamma) d\gamma dx.$$

- From a quantum physics perspective, we would like  $P(x, \gamma)$  to behave as a true probability distribution.
- The expected value of an observable  $G$  is given by

$$\langle Gf, f \rangle = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} g(x, \gamma) P(x, \gamma) d\gamma dx$$

where the operator  $G$  is given as the function  $g$  of the position and momentum operators, i.e.,

$$G = g(X, (-1/2\pi i)d/dx), Xf(x) = xf(x).$$

- These considerations led to the invention of the Wigner distribution.

## Definition (Wigner Distribution)

Given  $f \in L^2(\mathbb{R}^d)$ , the *Wigner distribution* of  $f$ , denoted  $Wf(x, \gamma)$  is defined by

$$Wf(x, \gamma) = \int_{\mathbb{R}^d} f\left(x + \frac{t}{2}\right) \overline{f\left(x - \frac{t}{2}\right)} e^{-2\pi i(\gamma \cdot t)} dt.$$

Given  $f, g \in L^2(\mathbb{R}^d)$ , the *cross-Wigner distribution* is given by

$$W(f, g)(x, \gamma) = \int_{\mathbb{R}^d} f\left(x + \frac{t}{2}\right) \overline{g\left(x - \frac{t}{2}\right)} e^{-2\pi i(\gamma \cdot t)} dt.$$

- Introduced by E. Wigner in 1932 in the context of quantum mechanical measurements.
- Starting with J. Ville in 1948 it has become popular and useful in signal analysis.

# Basic properties of $W(f, g)$

## Theorem

The following hold.

- If  $f, g \in L^2(\mathbb{R}^d)$ , then  $W(f, g) \in L^\infty \cap L^2(\mathbb{R}^{2d})$  is uniformly continuous on  $\mathbb{R}^{2d}$ , and  $Wf$  is real-valued.
- If  $f_1, f_2, g_1, g_2 \in L^2(\mathbb{R}^d)$ , Moyal's formula holds, i.e.

$$\langle W(f_1, g_1), W(f_2, g_2) \rangle_{L^2(\mathbb{R}^{2d})} = \langle f_1, f_2 \rangle_{L^2(\mathbb{R}^d)} \overline{\langle g_1, g_2 \rangle_{L^2(\mathbb{R}^d)}}.$$

- If  $f, \hat{f} \in L^1 \cap L^2(\mathbb{R}^d)$  then  $Wf(x, \omega)$  satisfies the marginals, i.e.
  - $\int_{-\infty}^{\infty} Wf(x, \gamma) d\omega = |f(x)|^2,$
  - $\int_{-\infty}^{\infty} Wf(x, \gamma) dt = |\hat{f}(\gamma)|^2,$

- If  $f(0) \neq 0$  then

$$f(x) = \frac{1}{f(0)} \int_{\mathbb{R}^d} Wf\left(\frac{x}{2}, t\right) e^{2\pi i \langle t, x \rangle} dt.$$

- This follows from the fact that the coordinate change

$$(x, y) \mapsto \left(x + \frac{y}{2}, x - \frac{y}{2}\right)$$

is inverted by

$$(x, y) \mapsto \left(\frac{x + y}{2}, x - y\right)$$

and then applying Fourier inversion.



# Positivity of $Wf$ .

- In order to be a true energy density function, we would like to say that  $Wf(x, \gamma) \geq 0$  for all  $f \in L^2$  and  $(x, \gamma) \in \mathbb{R}^{2d}$ . This only holds for functions of a special form, the *generalized gaussians*.
- However, positivity for  $Wf$  can be achieved by taking averages, i.e.,  $Wf * \sigma$ . This fact can be interpreted as a form of the uncertainty principle.

## Theorem

For  $a, b > 0$  let

$$\sigma_{a,b}(x, \gamma) = e^{-2\pi\left(\frac{x^2}{a} + \frac{\gamma^2}{b}\right)}.$$

Then

- $Wf * \sigma_{a,b} > 0$  for all  $f \in L^2(\mathbb{R}^d)$  if and only if  $ab > 1$ .
- $Wf * \sigma_{a,b} \geq 0$  for all  $f \in L^2(\mathbb{R}^d)$  if and only if  $ab \geq 1$ .
- This can be interpreted as saying that the values of  $Wf(x, \gamma)$  at a particular point  $(x, \gamma)$  or even in a set  $U \subseteq \mathbb{R}^{2d}$  of small area do not have a clear physical meaning.

# The Short-Time Fourier Transform

- From the late 1940s, communication engineers took a different (but related) approach to understanding the joint time-frequency content of a signal, the *spectrogram*.

## Definition

Given  $a, b \in \mathbb{R}^d$ , we define for  $f \in L^2(\mathbb{R}^d)$  the

- *translation operator*, or *time shift operator*,  
 $T_a f(x) = f(x - a)$ , and the
- *modulation operator*, or *frequency shift operator*,  
 $M_b f(x) = e^{2\pi i(b \cdot x)} f(x)$ .

Given  $g \in L^2(\mathbb{R}^d)$ , we define the *short-time Fourier transform (STFT)* on  $L^2(\mathbb{R}^d)$  by

$$V_g f(x, \gamma) = \int_{\mathbb{R}^d} f(t) \overline{g(t - x)} e^{-2\pi i(t \cdot \gamma)} dt = \langle f, M_\gamma T_x g \rangle.$$

- $V_g$  is also referred to as the *voice transform* or the *windowed Fourier transform* with  $g$  being the *window function*.
- $V_g f(x, \gamma) = \widehat{(f)(T_x \bar{g})}(\gamma)$ .
- For fixed window  $g$ , the function  $|V_g f(x, \gamma)|^2$  is called the *spectrogram* of  $f$ .

# Properties of the STFT

- For  $f, g \in L^2$ ,  $V_g f(x, \gamma)$  is uniformly continuous on  $\mathbb{R}^{2d}$  and bounded with

$$\|V_g f\|_\infty \leq \|f\|_2 \|g\|_2.$$

- For fixed  $g \in L^2$ ,  $V_g: L^2(\mathbb{R}^d) \rightarrow L^2(\mathbb{R}^{2d})$  and moreover

$$\|V_g f\|_{L^2(\mathbb{R}^{2d})} = \|f\|_{L^2(\mathbb{R}^d)} \|g\|_{L^2(\mathbb{R}^d)}$$

so that if  $\|g\| = 1$ ,  $V_g$  is an isometry.

- Moyal's formula is satisfied, i.e., if  $f_1, f_2, g_1, g_2 \in L^2(\mathbb{R}^d)$  then

$$\langle V_{g_1}(f_1), V_{g_2}(f_2) \rangle_{L^2(\mathbb{R}^{2d})} = \langle f_1, f_2 \rangle_{L^2(\mathbb{R}^d)} \overline{\langle g_1, g_2 \rangle_{L^2(\mathbb{R}^d)}}.$$

- $V_g$  is invertible with

$$f(t) = \frac{1}{\|g\|_2} \iint_{\mathbb{R}^{2d}} V_g f(x, \gamma) M_\gamma T_x g(t) d\gamma dt.$$

# The STFT as a TF Distribution

- Many of the fundamental properties of the Wigner distribution follow from those of the STFT by observing that

$$W(f, g^-)(x, \gamma) = 2^d e^{4\pi i(x \cdot \gamma)} V_g f(2x, 2\gamma)$$

where  $g^-(x) = g(-x)$ .

- Note that although

$$Wf(x, \gamma) = 2^d e^{4\pi i(x \cdot \gamma)} V_{f^-} f(2x, 2\gamma),$$

the time-frequency distribution  $V_{f^-} f(x, \gamma)$  does not satisfy the marginals.

# Radar Ambiguity Function

## Definition

The (*radar*) *ambiguity function* of  $f \in L^2(\mathbb{R}^d)$  is given by

$$Af(x, \gamma) = \int_{\mathbb{R}^d} f\left(t + \frac{x}{2}\right) \overline{f\left(t - \frac{x}{2}\right)} e^{-2\pi i(\gamma \cdot t)} dt.$$

and given  $g \in L^2(\mathbb{R}^d)$ , the *cross-ambiguity function* by

$$A(f, g)(x, \gamma) = \int_{\mathbb{R}^d} f\left(t + \frac{x}{2}\right) \overline{g\left(t - \frac{x}{2}\right)} e^{-2\pi i(\gamma \cdot t)} dt.$$

- The radar ambiguity function is due to Woodward, 1953.

# Why “radar”?

- Suppose that a test signal  $f(t)$  is sent from a radar apparatus a distance  $r$  from the source and moving at velocity  $v$  from the source.
- The echo of  $f$  is received with time-shift  $\Delta t$  (proportional to the range) and Doppler shift  $\Delta\gamma$  (proportional to the velocity), i.e.

$$e(t) = M_{\Delta\gamma} T_{\Delta t} f.$$

- The echo is compared to arbitrary time-frequency shifts of  $f$ , and we compute

$$|\langle \mathbf{e}, M_{\omega} T_x f \rangle| = |A f(x - \Delta t, \omega - \Delta\gamma)|$$

- Maximizing this quantity over all  $(x, \omega)$  gives estimates for  $(\Delta t, \Delta\gamma)$ .



# Uncertainty Principles

- The *uncertainty principle* can be loosely formulated as follows.

*A function and its Fourier transform cannot simultaneously be concentrated on small sets.*

- An alternate but related formulation says the following.

*The time-frequency distribution of any function cannot be concentrated on a small set in the time-frequency plane. This is independent of which time-frequency distribution is used.*

# Classical Uncertainty Principle

## Theorem

Let  $f \in L^2(\mathbb{R})$  and let  $a, b \in \mathbb{R}$  be given. Then

$$\left( \int_{\mathbb{R}} (x-a)^2 |f(x)|^2 dx \right)^{1/2} \left( \int_{\mathbb{R}} (\gamma-b)^2 |\hat{f}(\gamma)|^2 d\gamma \right)^{1/2} \geq \frac{1}{4\pi} \|f\|_2^2.$$

Equality holds if and only if  $f(x)$  is a time-frequency shift of a gaussian  $\varphi(x) = e^{-\pi x^2/c}$ .

- If  $\|f\|_2 = 1$  then  $|f(x)|^2$  and  $|\hat{f}(\gamma)|^2$  can be thought of as probability distributions, and the quantities

$$\min_{a \in \mathbb{R}} \left( \int_{\mathbb{R}} (x-a)^2 |f(x)|^2 dx \right)^{1/2}, \quad \min_{\gamma \in \mathbb{R}} \left( \int_{\mathbb{R}} (\gamma-b)^2 |\hat{f}(\gamma)|^2 d\gamma \right)^{1/2}$$

as their standard deviations and hence as a measure of concentration of  $|f|^2$  and  $|\hat{f}|^2$  about their means.

## Lemma

Let  $A$  and  $B$  be self-adjoint operators on a Hilbert space. Then

$$\|(A - a)f\| \|(B - b)f\| \geq \frac{1}{2} |\langle [A, B]f, f \rangle|$$

for all  $a, b \in \mathbb{R}$ , and  $f$  in the domain of  $AB$  and  $BA$ , where  $[A, B] = AB - BA$  is the commutator of  $A$  and  $B$ .

## Proof:

$$\begin{aligned} \langle [A, B]f, f \rangle &= \langle [(A - a), (B - b)]f, f \rangle \\ &= \langle (B - b)f, (A - a)f \rangle - \langle (A - a)f, (B - b)f \rangle \\ &= 2i \operatorname{Im} \langle (B - b)f, (A - a)f \rangle. \end{aligned}$$

Applying Cauchy-Schwarz yields the result.

Apply the above Lemma to the position and momentum operators,

$$Xf(x) = x f(x), \text{ and } Pf(x) = \frac{1}{2\pi i} f'(x)$$

which are self-adjoint and whose commutator is

$$[X, P]f = \frac{1}{2\pi i} (xf'(x) - (xf)'(x)) = -\frac{1}{2\pi i} f(x).$$

Finally observe that

$$\|(X - a)f\|_2 = \left( \int_{\mathbb{R}^d} (x - a)^2 |f(x)|^2 dx \right)^{1/2}$$

and

$$\|(P - b)f\|_2 = \|(\widehat{P - b})f\|_2 = \left( \int_{\mathbb{R}^d} (\gamma - b)^2 |\widehat{f}(\gamma)|^2 d\gamma \right)^{1/2}$$

# Uncertainty Principle of Donoho and Stark

Theorem (D. Donoho-P. Stark, 1989)

Suppose that for some sets  $T, \Omega \subseteq \mathbb{R}^d$ ,  $f \in L^2(\mathbb{R}^d)$  satisfies

$$\left( \int_{T^c} |f(x)|^2 dx \right)^{1/2} < \epsilon_T \|f\|_2, \text{ and } \left( \int_{\Omega^c} |\hat{f}(\gamma)|^2 d\gamma \right)^{1/2} < \epsilon_\Omega \|\hat{f}\|_2.$$

Then

$$|T||\Omega| \geq (1 - \epsilon_T - \epsilon_\Omega)^2.$$

- We say that under the hypotheses of the theorem  $f$  and  $\hat{f}$  are  $\epsilon_T$  (resp.  $\epsilon_\Omega$ ) *concentrated* on  $T$  (resp.  $\Omega$ ).
- The theorem can be interpreted as saying that any function must essentially occupy an area of at least one in time-frequency space.

# Uncertainty Principle for the STFT

Theorem (Lieb 1990, Gröchenig 2001)

Let  $\|f\|_2 = \|g\|_2 = 1$  and suppose that for some  $\epsilon > 0$  and  $U \in \mathbb{R}^{2d}$ ,

$$\iint_U |V_g f(x, \gamma)|^2 dx d\gamma \geq 1 - \epsilon.$$

Then  $|U| \geq 2^d(1 - \epsilon)^2$ .